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TECHNICAL NOTE 2374

EFFECT OF INITIAL MIXTURE TEMPERATURE ON
FLAME SPEED OF METHANE-AIR, PROPANE-AIR
AND ETHYLENE-AIR MIXTURES

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SUMMARY

Flame speeds based on the inner edge of the shadow cast by the laminar Bunsen cone were determined for methane-air and ethylene-air mixtures as functions of composition at initial mixture temperatures ranging from 34° to 344° C. The data showed that flame speed increased with temperature at an increasing rate, as did similar data reported earlier for propane and air. The flame speeds were found to be affected by temperature on the basis of percentage in the decreasing order, methane, propane, and ethylene. The flame speed u for none of these fuels was exactly proportional to a power of initial temperature T_0 ; however, approximate equations for flame speed in centimeters per second (based on inner edge of shadow) for use in the temperature range 34° to 344° C were determined to be as follows:

For methane,

$$u = 0.000780 T_0^{1.89}$$

for propane,

$$u = 0.00246 T_0^{1.72}$$

and for ethylene,

$$u = 0.00971 T_0^{1.565}$$

Part of the data were recomputed on the basis of the surface indicated by the outer edge of the shadow cast by the flame cone. For a given fuel, the ratio of outer-edge value to inner-edge value was constant with respect to both temperature and composition, within the accuracy of the outer-edge measurements. Flame speeds calculated from the preceding empirical equations could therefore be converted to outer-edge values by multiplying by the constants 0.900, 0.885, and 0.849 for

methane, propane, and ethylene, respectively. Outer-edge values (room temperature) were close to values obtained from the visible Bunsen cone and from tube measurements.

The observed effect of temperature on flame speed for each of the fuels was reasonably well predicted by either the thermal theory as presented by Semenov or the square-root law of Tanford and Pease. The importance of active radicals in flame propagation was indicated by a simple linear relation between maximum flame speed and equilibrium radical concentrations for all three fuels. Equally good correlations resulted from using either hydrogen-atom concentration alone or a summation of effective relative concentrations of hydrogen atoms, hydroxyl radicals, and oxygen atoms and using either flame temperatures based on a sodium D-line measurement for a room-temperature mixture or adiabatic flame temperatures.

INTRODUCTION

The flame speed, or normal burning velocity, of a fuel-air mixture is a fundamental property governing flame propagation, which is one of the several major processes occurring in combustion equipment for flight propulsion. Inasmuch as plots of the performance data of aircraft combustors show that the combustion efficiency is related to the combustor-inlet temperature (for example, references 1 and 2), knowledge of how temperature affects the fundamental combustion properties of the fuel-air mixture is desirable. Consequently, as part of the combustion research program at the NACA Lewis laboratory, an investigation was initiated to study the effect of initial mixture temperature on flame propagation. The effect of temperature on flame speeds and stability limits of propane-air flames is reported in reference 3. An extension of the flame-speed phase of this investigation to include the effect of temperature on methane-air and ethylene-air flames is reported herein.

Previous work (reference 3) showed that the flame speed of propane-air mixtures increased with temperature at an increasing rate. It was shown that the thermal theory of flame propagation, as presented by Semenov in reference 4 for bimolecular reactions, predicted the effect of initial mixture temperature on the maximum flame speed of propane-air mixtures within approximately 5 percent. It was further shown that the square-root law of Tanford and Pease (references 5 and 6), which is based on the diffusion of active particles, predicted relative increases in maximum flame speed that were as much as 35 percent lower than the experimental results. Subsequent reevaluation of the temperature dependence predicted by the square-root law has shown, however, that the relative increases predicted by this theory for propane are within 14 percent of the experimental results.

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Inasmuch as the activation energy of the oxidation process is an important factor in the thermal-theory equations, it was decided to obtain flame-speed - temperature data for a gaseous fuel having an activation energy appreciably different from that of propane. Methane was selected because it has an activation energy of 52 kilocalories per gram-mole (from unpublished data obtained by A. D. Walsh, Leeds University, England), compared with 38 kilocalories per gram-mole for propane (reference 7, p. 437).

It was also thought that data for ethylene, which has shown anomalous behavior in other correlations (reference 8) when compared with other hydrocarbons, would be of value. The activation energy of ethylene, which is approximately 40 kilocalories per gram-mole (reference 9), does not differ much from that of propane.

The present report contains the results of a study of the effect of initial mixture temperature over the range of 34° to 344° C on the laminar flame speeds of methane-air and ethylene-air mixtures. The flame speeds were computed from measurements based on the inner edge of the shadow cast by the Bunsen cone; the shadowgraphs were obtained with a parallel-beam shadowgraph system. The flame-speed values thus obtained are compared with values obtained by other methods.

Methane, propane, and ethylene are compared on the basis of relative increase in flame speed with initial temperature. A comparison is also made among the experimental data for each fuel and the relative effects of temperature on flame speed predicted by the thermal and diffusion theories.

SYMBOLS

The following symbols are used in this report:

- A longitudinal cross-sectional area of cone, (sq cm)
 D_i diffusion coefficient of i^{th} radical at initial temperature
 $D_{i,r}$ relative diffusion coefficient of given active radical with respect to other radicals
 D_m diffusion coefficient of radical at mean combustion-zone temperature
E activation energy, (cal/gram-mole)

h	height of cone, (cm)
k	constant for given fuel
l	slant height or length of generating curve, (cm)
n	constant exponent for given fuel
p_i	mole fraction or partial pressure of i^{th} radical in burned gas
R	gas constant, (cal/(gram-mole)(°K))
S	lateral surface area, (sq cm)
$T_{\text{av}} = (T_0 + T_f)/2$	
T_f	flame temperature, (°K)
T_0	initial mixture temperature, (°K)
u	flame speed, (cm/sec)

EXPERIMENTAL PROCEDURE

Apparatus. - The apparatus and the experimental technique used for the measurement of flame speed were essentially the same as reported in reference 3. The apparatus is diagrammatically illustrated in figure 1. In brief, the fuel and the air were metered, mixed, electrically preheated, and burned above a vertical brass tube. For methane, a 120-centimeter length of 15.7-millimeter-inside-diameter tubing was used; whereas, for ethylene, a 120-centimeter length of 6.3-millimeter-inside-diameter tubing was used. The smaller tube diameter for ethylene was necessary because of the higher flashback velocity limit of ethylene, which results from its higher flame speed. Temperatures were measured with iron-constantan thermocouples. The gas temperature at the burner port was periodically checked by locating an aspirating thermocouple (reference 10) over the center of the port. The gas temperatures at the burner-tube inlet (thermocouple "O") and at the port (thermocouple "C") were maintained within 10° C of each other. The parallel-beam shadowgraph system is schematically shown above the burner port in figure 1.

One modification to the original apparatus consisted in fastening a small resistance heater to the bottom of the burner-lip collar to permit close control of the lip temperature. In the propane-air investigation, the lip temperature was established by the heating due to the flame itself.

Fuels. - The minimum purities claimed by the suppliers of the methane and ethylene were 99.0 and 99.5 percent, respectively. The propane (reference 3) had a minimum purity of 95.0 percent, the major impurities being ethane and isobutane. Laboratory service air containing approximately 0.3-percent water by weight was used.

Computation of flame speed. - Flame speeds were determined from the shadowgraphs by the total-area method, wherein the average normal flame speed is equal to the volume rate of flow of the unburned mixture divided by the surface area of the cone formed by the combustion zone. This surface area S was determined by assuming that the flame surface can be approximated by the relation for conical surfaces of revolution; thus,

$$S = \pi A l/h \quad (1)$$

A typical shadowgraph is shown in figure 2. The longitudinal cross-sectional cone area A was determined by measuring a magnified image of the shadow cone with a planimeter; the magnified image was obtained by tracing the inner edge of the shadow from the projected image of the shadowgraph. Because of inaccuracies in tracing and using a planimeter, a scatter in the computed points resulted, especially for the tall, slender ethylene flames. The precision at the higher temperatures was 4 percent for ethylene. This scatter was diminished by: (1) determining the average value of A/h for all the computed points at a given temperature, (2) recomputing flame-speed values for one of the points on the basis of the average A/h ; and (3) computing the rest of the flame-speed values from the already corrected point by a ratio of cone heights according to the equation $u_2 = u_1 h_1/h_2$. This procedure essentially amounts to the method used in reference 11.

RESULTS AND DISCUSSION

In figure 3, flame speed is plotted against percentage of fuel by volume at each of six initial mixture temperatures ranging from 34° to 344° C. It is assumed that no preflame reaction has occurred during the short preheating period (less than 2 sec) at these temperatures

(reference 7, pp. 406, 411). The stream-flow Reynolds number was approximately 1500 for methane-air mixtures and 2000 for ethylene-air mixtures. Each of the curves shows that the maximum flame speed for a given temperature occurs at a mixture composition richer than stoichiometric. These flame-speed maximums from figure 3 are plotted against temperature in figure 4(a) to show the effect of temperature on the maximum flame speed. The curve from reference 3 for the propane-air mixture is included for comparison. For further comparison, a log-log plot of maximum flame speed against temperature for the three fuels is presented in figure 4(b). None of the curves in figure 4(b) may be exactly represented by a straight line. For an approximate relation, however, an equation of the type

$$u = k T_0^n$$

may be used. The equations, found by the least-squares method from the points plotted in figure 4(b) (based on the inner edge of the shadow) for the three fuels, are:

For methane,

$$u = 0.000780 T_0^{1.89}$$

for propane,

$$u = 0.00246 T_0^{1.72}$$

and for ethylene,

$$u = 0.00971 T_0^{1.565}$$

These values should not be used for extrapolation to higher or lower temperatures until further experimental data are available.

In figure 5, the percentage increase in maximum flame speed, referred to the flame speed at 34° C, is plotted against temperature. (The value for propane at this temperature was obtained by interpolation from fig. 3 of reference 3.) The flame speeds of these three fuels are again seen to be affected by temperature in the decreasing order, methane, propane, and ethylene.

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Flame speeds based on outer edge of shadow. - As indicated in references 9 and 12, the maximum density gradient is more nearly represented by the outer edge of the shadow than by the inner edge, because of the manner in which the shadow of a flame cone is produced. If the locus of maximum density gradient is taken to correspond to the flame front, then absolute values of flame speed should be based on the outer edge of the shadow. In order to obtain correction factors that could be used to convert the flame-speed values plotted in the figures of the present report to values based on the outer edge of the shadow, a group of points was selected at each temperature for each fuel and was recomputed on the new basis. Each group consisted of one point representing the leanest mixture studied, one point near the composition for maximum flame speed, and one for the richest mixture studied. Plots of 18 points so selected for each fuel revealed no trends for the correction factor with respect to either mixture composition or temperature; the arithmetic mean value of the factors for all 18 points for a given fuel was therefore considered to be the best value over the ranges of composition and temperature studied. Inasmuch as all data points for one fuel would then be corrected by the same constant factor for that fuel, such correction would not change the relative effects of temperature on flame speed shown in figures of this report; that is, it would merely shift the ordinate log scale in figure 4(b) and compress the ordinate scales in the remaining figures. The empirical equations obtained from figure 4(b) could therefore be converted to give values based on the outer edge of the shadow by multiplying by the constants 0.900, 0.885, and 0.849 for methane, propane, and ethylene, respectively. Values based on the inside edge of the shadow are presented herein because the required measurements were more easily and precisely made on this basis and because the relative effect of temperature on flame speed was unaffected.

The arithmetic mean values of the correction factors for the three fuels are presented in table I. The maximum flame speed at 25° C was obtained for each fuel by extrapolation from figure 4(a) in order to compare the value with values obtained by other investigators who used shadow-cone, visible-cone, and tube measurements (references 3, 8, 9, and 13 to 16). The agreement among corrected shadow-cone, visible-cone, and tube measurements for methane and propane is good. The poorer agreement between various methods and investigators for ethylene is probably due, in part, to the greater uncertainties in measurements of the more extreme flame shapes encountered with ethylene in both Bunsen burner and tube methods.

Comparison of experimental data with relative values predicted by theoretical equations. - It is shown in reference 3 that, for the purpose of predicting the relative effect of temperature on flame speed, the thermal-theory equation of Semenov (reference 4), for a bimolecular reaction controlling, might be reduced to the expression

$$u \propto \sqrt{T_f^{4.9} T_0^2 \frac{e^{(-E/RT_f)}}{(T_f - T_0)^3}} \quad (2)$$

The reduction of the Semenov equation to expression (2) involves the assumption that the relations between the physical properties and the temperature for the mixture are reasonably near those for air.

In a similar manner, the thermal-theory equation for unimolecular reaction controlling (reference 4) may be reduced to the following expression:

$$u \propto \sqrt{T_f^{3.83} T_0^2 \frac{e^{(-E/RT_f)}}{(T_f - T_0)^2}} \quad (3)$$

The square-root law of Tanford and Pease in the form presented in reference 6 was also considered for the purpose of predicting the relative effect of temperature on flame speed in reference 3. The reduced expression presented in reference 3 is

$$u \propto \sqrt{\left(\sum_i p_i D_{i,r} \right) T_0^{1.67} T_{av}^{-1}} \quad (4)$$

where

$$\sum_i p_i D_{i,r} = 6.5 p_H + p_{OH} + p_O$$

In arriving at expression (4), it was assumed that the number of molecules per unit volume of gas at some mean temperature L , should be inversely proportional to T_{av} . In reference 5, however, Tanford and Pease apparently intended the mean temperature to be chosen nearer to and perhaps proportional to the flame temperature, for example, $0.7 T_f$. The diffusion coefficients were also assumed to be proportional to the square of the absolute temperature, rather than to the 1.67 power used in reference 3. Cancellations of temperature effects involving the two-power relation were made in the process of substituting the diffusion coefficient of radicals at the initial temperature D_i for the diffusion coefficient of radicals at the mean combustion-zone temperature D_m (reference 5); these cancellations resulted in the equation presented in reference 6. If, instead, the temperature dependence $D \propto T^{1.67}$ is applied to the D_m in the derivation (reference 5) rather than to D_i in the final equation as in reference 3, and if $L \propto T_f^{-1}$ is substituted, the revised expression is

$$u \propto \sqrt{\left(\sum_i p_i D_{i,r} \right) \frac{T_0^2}{T_f^{1.33}}} \quad (5)$$

Comparisons of the relative effects of initial temperature on flame speed as predicted by the reduced equations (2), (3), and (5) for methane, propane, and ethylene are presented in figures 6(a), 6(b), and 6(c), respectively. Each of the curves was plotted by computing flame speeds at various temperatures relative to the experimental value of maximum flame speed at room temperature. The experimental results are represented by the dashed curves.

The activation energies and the flame temperatures used in plotting the theoretical curves in figure 6 are presented in table II. The flame temperatures are based on sodium D-line measurements of flame temperatures of mixtures at room temperature (reference 17). The relations for the change in flame temperature ΔT_f with change in initial temperature ΔT_0 were determined from adiabatic flame-temperature calculations by the method of reference 18. Equilibrium radical concentrations needed for the square-root-law curves were calculated by the graphical method of reference 19 and are presented in table III. For each of the fuel-air mixtures (air containing 0.3-percent water by weight), the radical concentrations were calculated for a series of arbitrary temperatures, which included the range of flame temperatures

considered. Intermediate values of radical concentrations corresponding to specific flame temperatures, which in turn corresponded to specific initial temperatures, were obtained from table III by logarithmic interpolation. Although the radical concentrations for methane-air flames in table III were computed for a fuel concentration of 10.0 percent by volume, whereas the flame-speed maximums in figure 3 occurred at approximately 10.5 percent methane, it can be shown that the curves in figures 6 and 7 would be insensitive to this small change in concentration, because the flame speeds at a concentration of 10.0 percent are very near the maximum values.

The curves in figure 6 show that both the expressions based on the thermal theory and the expression based on a diffusion theory predict the relative effect of initial mixture temperature on maximum flame speed within 20 percent. For methane and propane, the thermal-theory curves are in somewhat better agreement with the experimental results than are the square-root-law curves; whereas for ethylene, the square-root-law curve is in as good, or better, agreement with the experimental results.

It must be emphasized that even qualitative observations regarding the ability of one or the other of the theories to predict the effect of initial temperature on flame speed are subject to the assumptions made in deriving both the original equations and the reduced forms and to the choices of values for E , T_f , and ΔT_f . The effects of using different values of E , T_f , and ΔT_f for propane flames is discussed in reference 3. For example, it could be shown that using $E = 25$ kilocalories per gram-mole, $T_f = 2390^\circ \text{K}$ at $T_0 = 305^\circ \text{K}$, and $\Delta T_f = 0.45 \Delta T_0$ for propane would give predicted increases that were approximately 17 percent low by the thermal-theory bimolecular equation and 17 percent low by the square-root law. It is assumed herein that the variation of B_i , a term that appears in the square-root law to allow for radical recombination, with initial temperature is relatively small and would not appreciably affect the curves in figure 6. The temperature dependence of k_i , the rate constant which appears in the square-root law for interaction between the i^{th} radical and the combustible material, has been neglected. Constant values of k_i were used for hydrogen, carbon monoxide, and methane flames in reference 6, where the variations in flame temperature are of the same order of magnitude as those considered herein. If the temperature dependence of k_i is represented by $e^{(-E/RT_f)}$, the temperature effect on flame speed due to k_i will be small, provided that the activation energy for the radical-hydrocarbon reaction is small. For example, using an activation

energy of 7 kilocalories per gram-mole (reference 6), the relative increases in flame speed predicted by the square-root law for methane, propane, and ethylene would be raised approximately 8 percent at an initial temperature of 344° C and less at lower initial temperatures.

The relative merits of the theories could be better assessed if reliable data for the effect of pressure on flame speed were available. The thermal-theory equations presented in reference 4 predict that for a unimolecular reaction, the flame speed should be inversely proportional to the square root of the absolute pressure; whereas, for a bimolecular reaction, the flame speed should be independent of pressure. The square-root law predicts that flame speed should vary approximately inversely as the fourth root of the pressure (reference 6). If one of the theories could be shown to be the more consistent in predicting both pressure and temperature effects, it would be strengthened considerably. Although some pressure data exist with which this precept might be tested, conflicting trends for the variation of flame speed with pressure for a given fuel have been reported by different investigators, or by the same investigator using different experimental methods, which indicates that some of these data were influenced by the experimental apparatus. For example, in reference 20 the flame speed of ethylene measured by a Bunsen burner method was found to be inversely proportional to the fourth root of pressure; whereas, according to reference 21 the flame speed of ethylene measured by a soap-bubble method was found by the same investigator to be independent of pressure.

Correlation of flame speeds with active-radical concentrations. -

A correlation between flame speed and $\sum_i p_i D_{i,r}$ for ethylene-air mixtures of various compositions at room temperature is found in reference 9. In figure 7(a), the maximum flame speeds from figure 4 are plotted against $\sum_i p_i D_{i,r}$ computed at temperatures determined from table II; in figure 7(b), maximum flame speeds are plotted against computed hydrogen-atom concentrations at these same temperatures; and in figure 7(c) maximum flame speeds are plotted against hydrogen-atom concentrations at theoretical flame temperatures obtained from reference 18. In all cases, straight-line correlations are obtained. The equations of these straight lines are given on the figures. The fact that the methane and propane points fall on a common line in figure 7(a) and 7(b) may be a fortuitous result of the choices of measured flame temperatures, inasmuch as they do not fall on a common line in figure 7(c), where theoretical flame temperatures were used. The fact

that the ethylene line is different in all cases was not too unexpected in view of the anomalous behavior of ethylene in a similar correlation in reference 8, in which flame speed was plotted against $\sum_i p_i D_{i,r}$ for a large number of fuels at room temperature.

The curves of figure 7 indicate that, at least for the three fuels studied, the maximum flame speed of a gaseous fuel may be accurately determined over a considerable range of initial mixture temperature when the experimental values at two temperatures are known and sufficient data exist to compute the adiabatic flame temperatures, and hence, the equilibrium hydrogen-atom concentrations. It is recognized that the slopes of the correlation lines in figure 7 are much steeper than those shown in the correlations in references 8 and 9. No explanation of the differences in the slopes is offered at the present time. Further experiments and calculations are indicated, the results of which may give clues as to the controlling mechanism of flame propagation under various conditions.

SUMMARY OF RESULTS

An investigation of the effect of initial mixture temperature on the laminar flame speeds of methane-air and ethylene-air mixtures, together with earlier work with propane-air mixtures, gave the following results:

1. The flame speeds of methane, propane, and ethylene increased with initial mixture temperature at an increasing rate. The percentage change in flame speeds was affected by temperature in the decreasing order, methane, propane, and ethylene. Although none of these gases showed maximum flame speed u to be exactly proportional to a power of temperature T , approximate equations for flame speed in centimeters per second (based on the inner edge of the shadow cast by the cone) for use in the temperature range 34° to 344° C were:

For methane,

$$u = 0.000780 T_0^{1.89}$$

for propane,

$$u = 0.00246 T_0^{1.72}$$

and for ethylene,

$$u = 0.00971 T_0^{1.565}$$

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2. Flame-speed values based on the outer edge of the shadow cast by the Bunsen cone (total-area method) were close to those based on both visible Bunsen cone measurements and tube measurements for methane and propane, but the variation in values obtained by the three methods was greater for ethylene. For each fuel, a factor was determined to allow conversion of values based on the inner edge of the shadow to values based on the outer edge of the shadow; the conversion factors were constant with respect to both temperature and composition, within the accuracy of the measurements based on the outer edge of the shadow. The empirical equations presented in result (1) could therefore be converted to give values based on the outer edge of the shadow by multiplying by the constants 0.900, 0.885, and 0.849 for methane, propane, and ethylene, respectively.

3. Both the thermal-theory equations presented by Semenov and the square-root law of Tanford and Pease (diffusion theory) may be used to predict within 20 percent or better the relative effect of initial temperature on flame speed for the gases studied. Of the two theories, the thermal theory gave somewhat better predictions for methane and propane, whereas the square-root law was as good, or better, for ethylene; this qualitative observation was subject, of course, to the assumptions made in applying the equations.

4. Very good straight-line correlations between maximum flame speed and either the hydrogen-atom concentration alone or the summation of the effective relative concentrations of hydrogen atoms, hydroxyl radicals, and oxygen atoms were found for the three fuels. Equally good correlations were obtained using either flame temperatures based on sodium D-line measurements or adiabatic flame temperatures. The slopes of the correlation lines were steeper than those reported elsewhere in investigations of the effects of mixture composition and fuel type on flame speed. Further experimental and theoretical investigation may show why the slopes were different and give a better insight into the mechanisms involved in flame propagation.

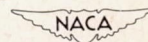
Lewis Flight Propulsion Laboratory,
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TABLE I - CORRECTION FACTORS FOR COMPUTING FLAME SPEEDS BASED ON
OUTER EDGE OF SHADOW AND COMPARISON OF 25° C VALUES WITH
VISIBLE BUNSEN-CONE AND TUBE MEASUREMENTS



Fuel	Correction factor		Flame speed at 25° C (cm/sec)				Flame speed at 340° C (cm/sec)	
	Mean value	Mean devi- ation	Inner edge	Corrected to outer edge	Visible cone	Tube method (a)	Inner edge (b)	Corrected to outer edge
Methane	0.900	0.017	b37.2	33.5	33.3	33.8	147	132
Propane	0.885	0.017	b44.9 d45.5 f46.6	40.8	c38.5 e41.5	39.0	158	140
Ethylene	0.849	0.027	b73.7 f74.6 h72.0	62.5	g60.0	68.3	229	195

^aBased on total area of flame traveling at uniform velocity inside pyrex tube (reference 8).

^bObtained from figure 4(a) by extrapolation or interpolation.

^cBased on total-area method, center of luminous zone (reference 3).

^dTotal-area value extrapolated to diametral axis of flame; value of 45.0 cm/sec obtained by stroboscopic method (reference 13).

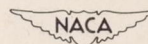
^eBased on area of cone frustum (reference 15).

^fBased on area of upper part of shadow cone only (reference 14).

^gBased on total-area method (reference 16).

^hMean value of number of measurements by angle method at different cone positions (reference 9).

TABLE II - ACTIVATION ENERGIES AND FLAME TEMPERATURES
USED FOR THEORETICAL CURVES



Fuel	E cal ($\frac{\text{gram-mole}}{\text{gram-mole}}$)	T _f (°K) (a)	ΔT _f (b)
Methane	^c 52,000	2150	0.54 ΔT ₀
Propane	^d 38,000	2200	0.50 ΔT ₀
Ethylene	^e 40,000	2250	0.475 ΔT ₀

^aFor T₀ = 307° K, based on sodium D-line measurement (reference 17).

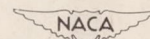
^bFrom adiabatic flame-temperature calculations (reference 18).

^cUnpublished data obtained by A. D. Walsh.

^dReference 7, p. 437.

^eReference 9.

TABLE III - EQUILIBRIUM RADICAL CONCENTRATIONS BY
GRAPHICAL METHOD^a



Fuel	Fuel (percent by volume)	T (°K)	Partial pressure, atm × 10 ³			$\sum_i p_i D_{i,r} \times 10^3$
			P _H	P _{OH}	P _O	
Methane	10.0	2100	0.274	0.811	0.0221	2.61
		2200	.494	1.76	.0836	5.05
		2300	.862	3.34	.262	9.20
		2400	1.47	5.98	.728	16.26
Propane	4.2	2200	0.430	1.65	0.088	4.53
		2300	.775	2.90	.250	8.18
		2400	1.35	5.20	.720	14.68
Ethylene	7.5	2200	0.665	0.913	0.0321	5.26
		2300	1.12	1.86	.113	9.25
		2400	1.83	3.39	.331	15.52
		2500	2.96	5.78	.857	25.88

^aComputed for air containing 0.3 percent water by weight (reference 19).

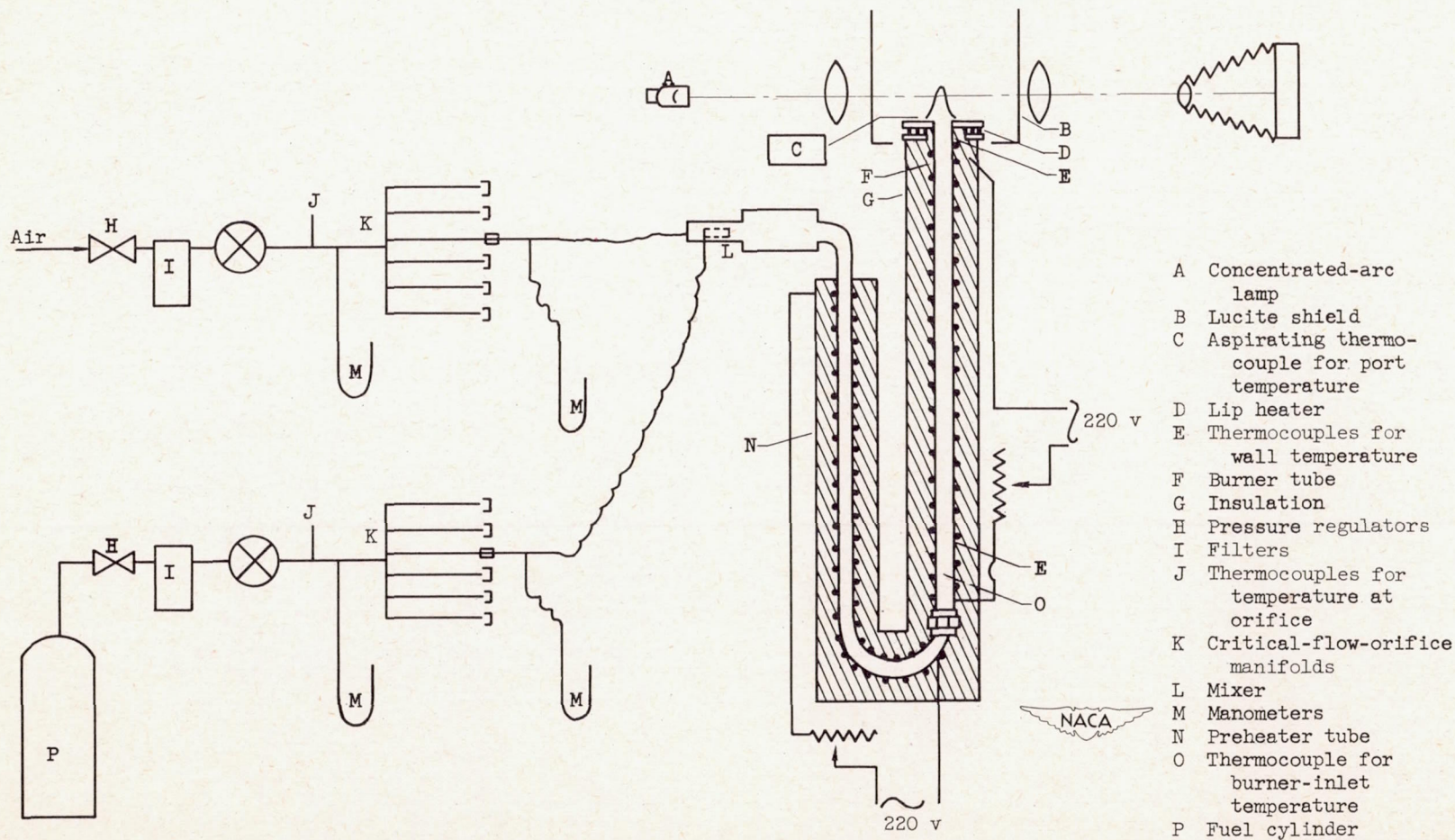


Figure 1. - Diagrammatic sketch of experimental apparatus.

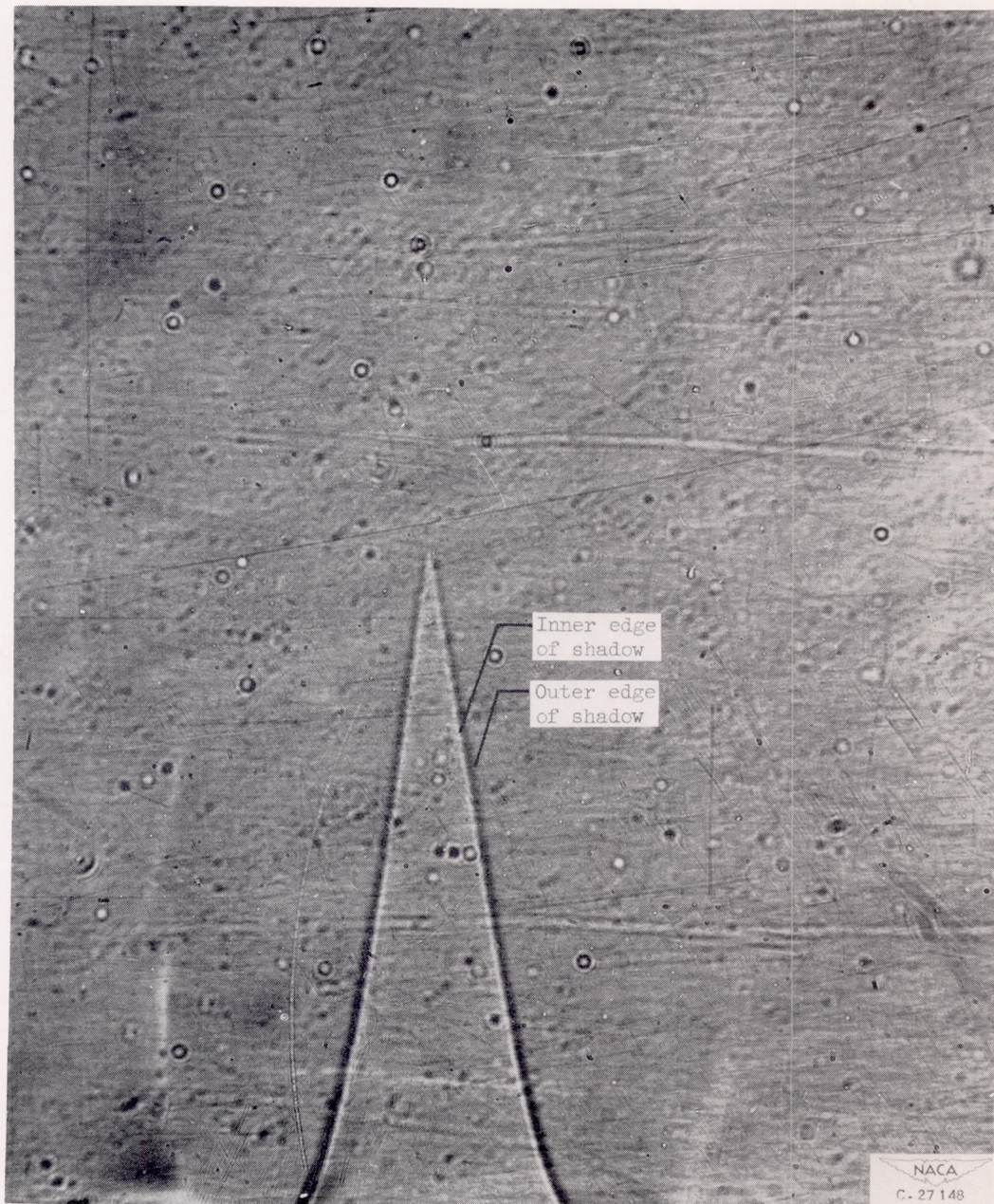
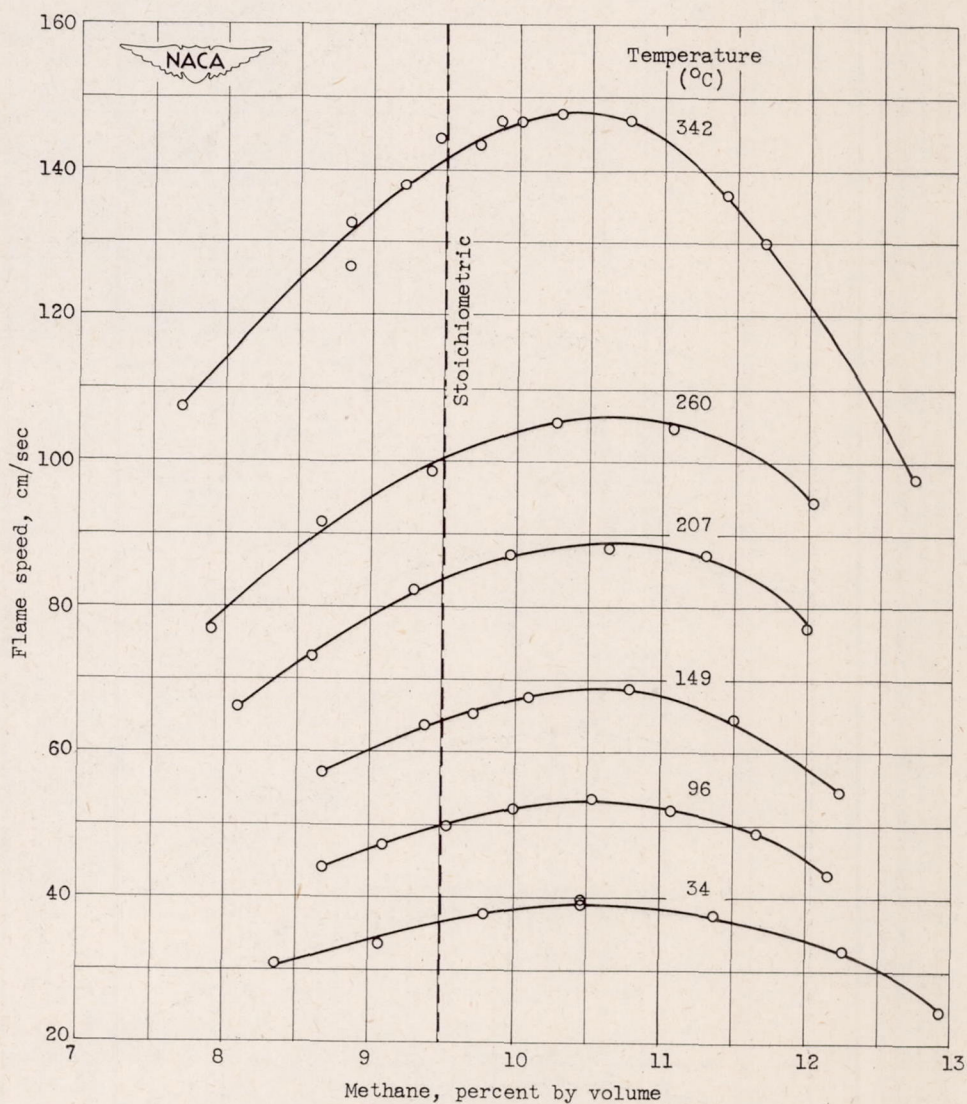
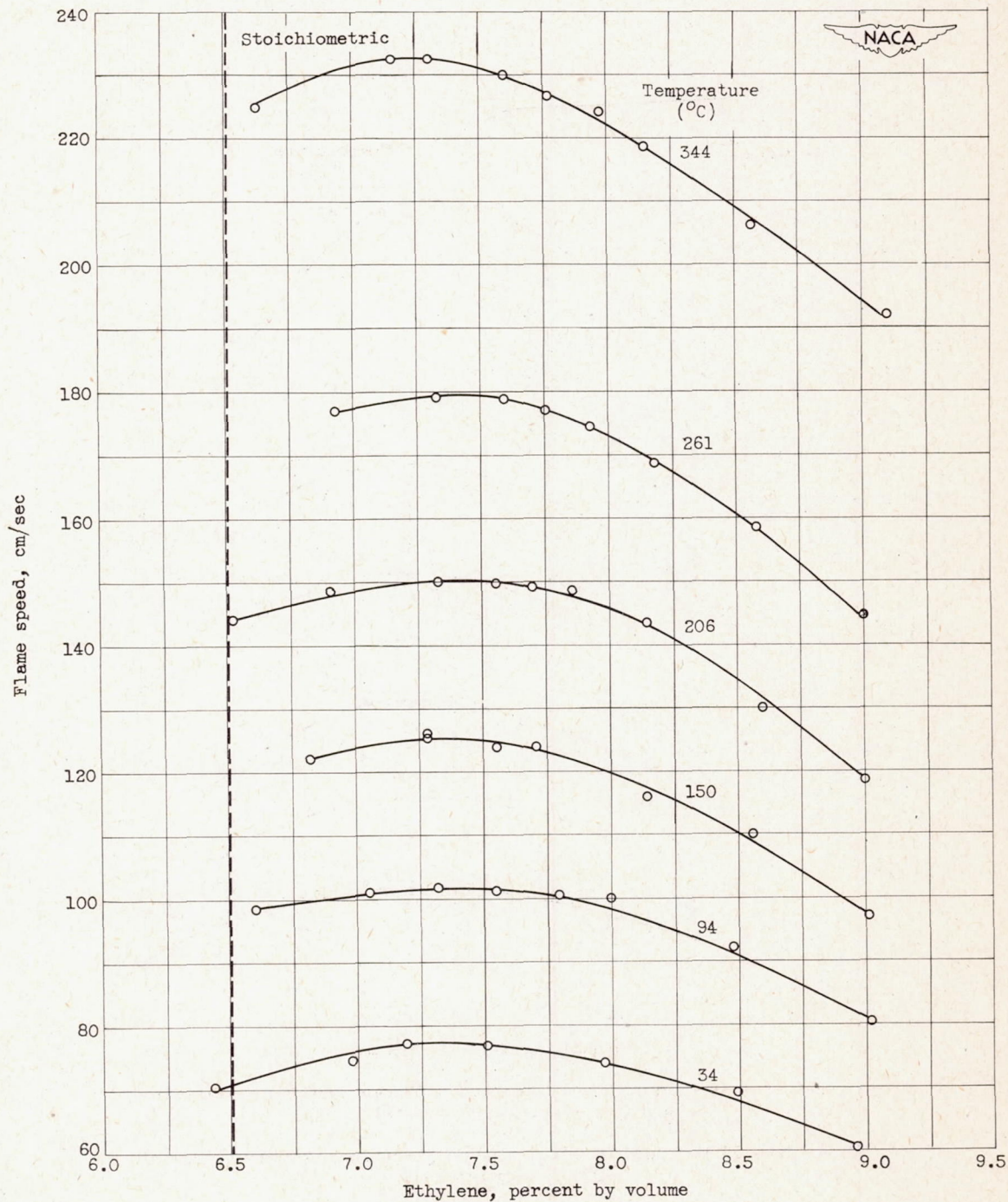


Figure 2. - Typical shadowgraph, showing inner and outer edges of shadow cast by Bunsen cone.



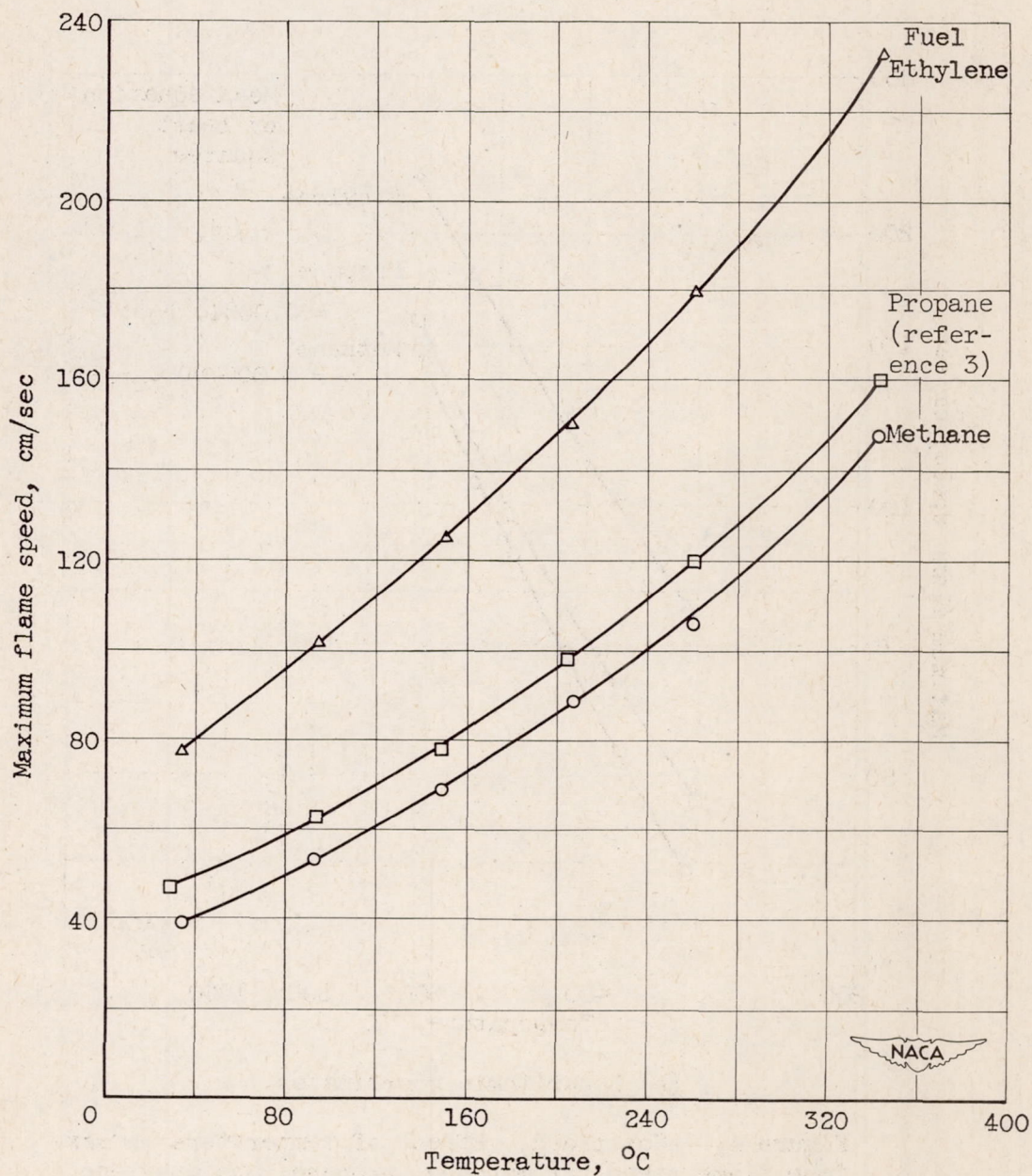
(a) Methane-air flames. Tube diameter, 1.57 centimeters; stream-flow Reynolds number, approximately 1500.

Figure 3. - Flame speed as a function of percentage fuel by volume at various temperatures.



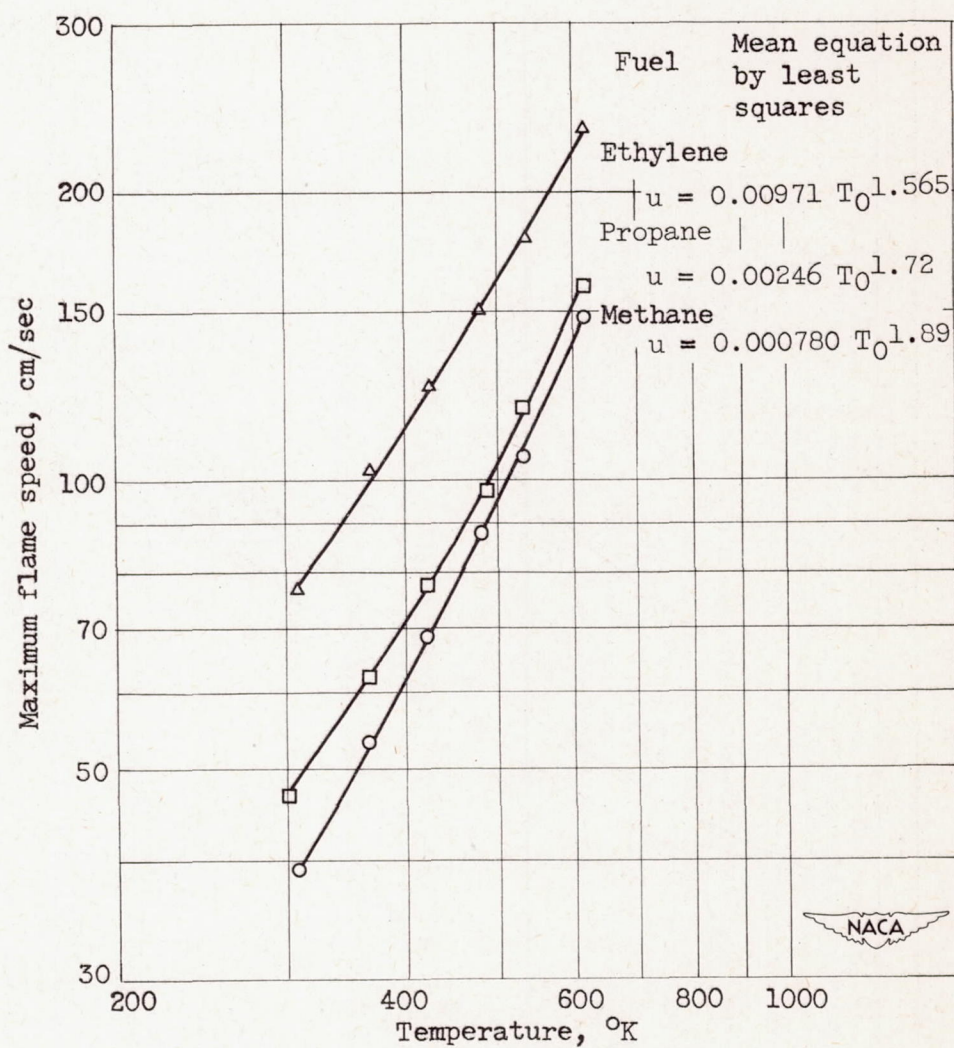
(b) Ethylene-air flames. Tube diameter, 0.631 centimeter; stream-flow Reynolds number, approximately 2000.

Figure 3. - Concluded. Flame speed as a function of percentage fuel by volume at various temperatures.



(a) Rectangular coordinates.

Figure 4. - Effect of temperature on maximum flame speed for methane, propane, and ethylene.



(b) Logarithmic coordinates.

Figure 4. - Concluded. Effect of temperature on maximum flame speed for methane, propane, and ethylene.

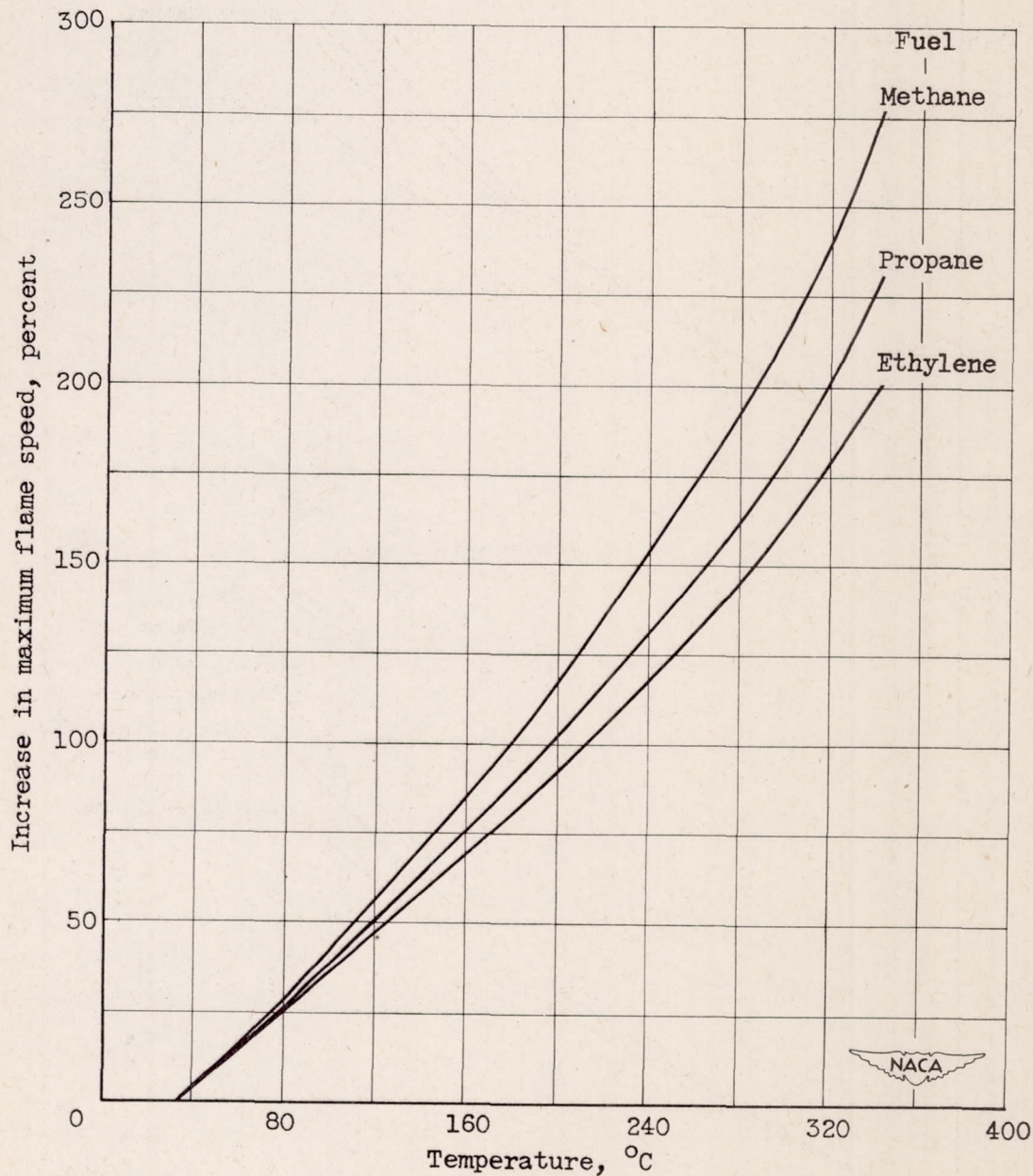
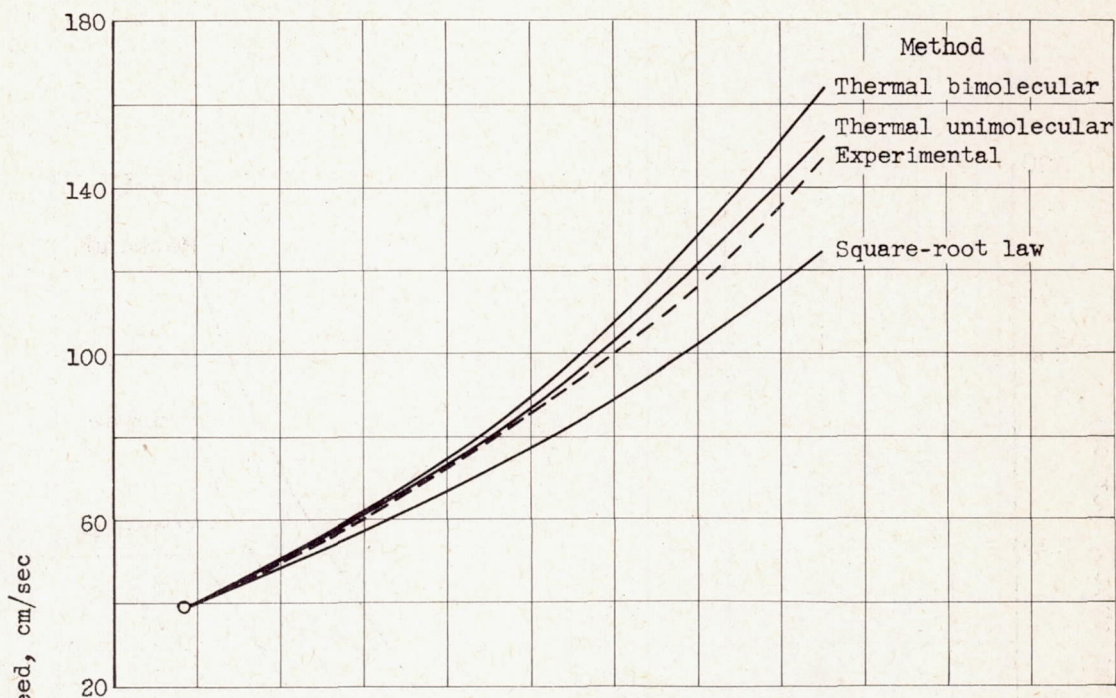
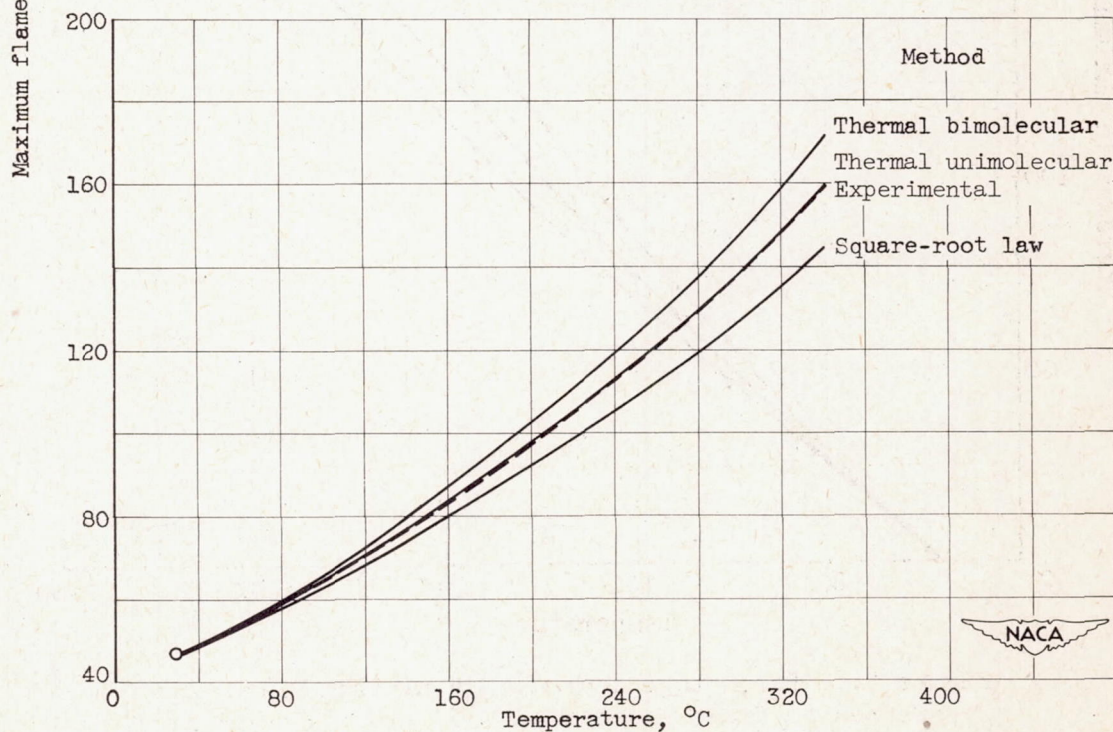


Figure 5. - Percentage increase of maximum flame speed, referred to flame speed at 34° C, with temperature for methane, propane, and ethylene.

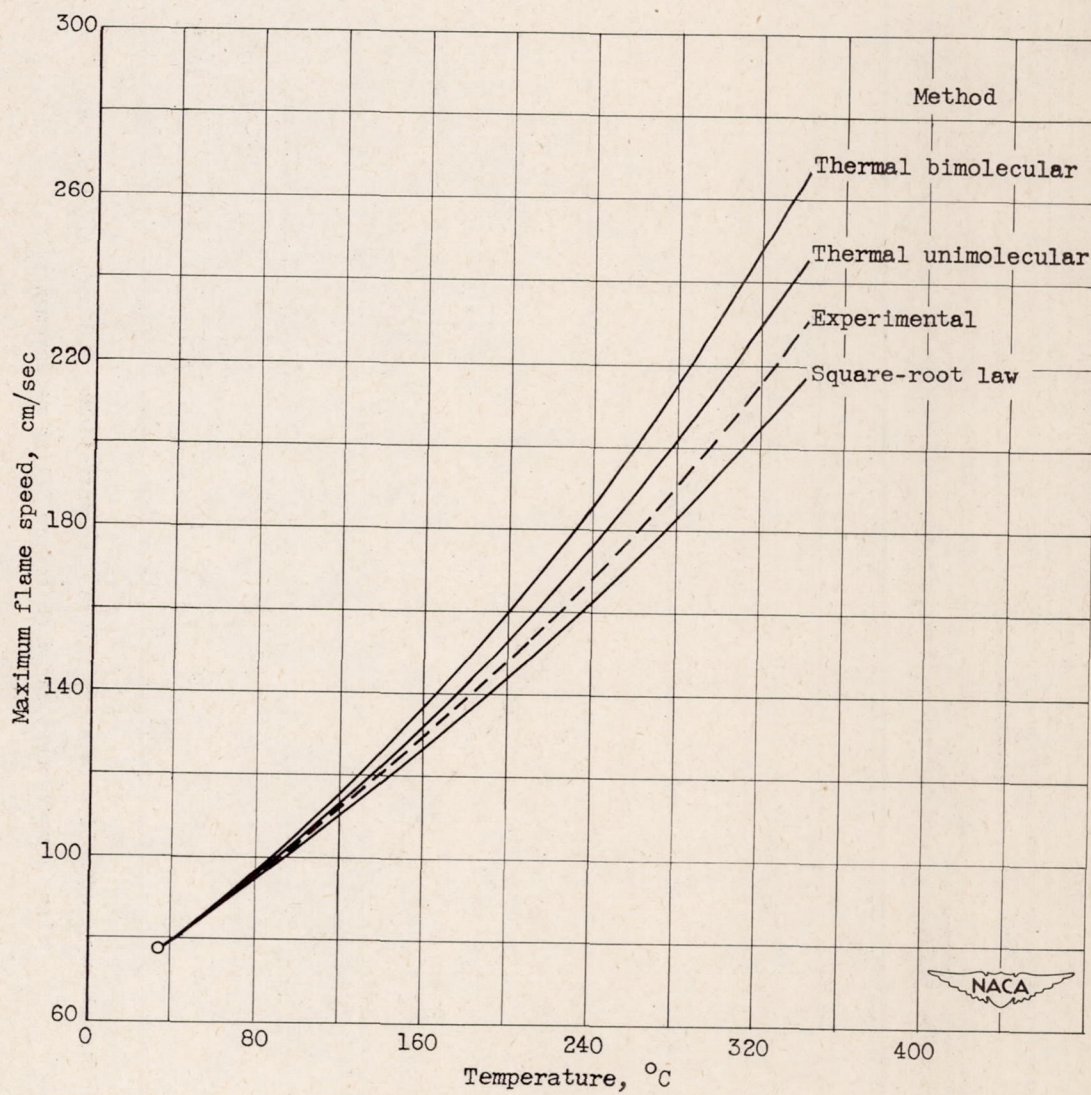


(a) Methane-air flames.



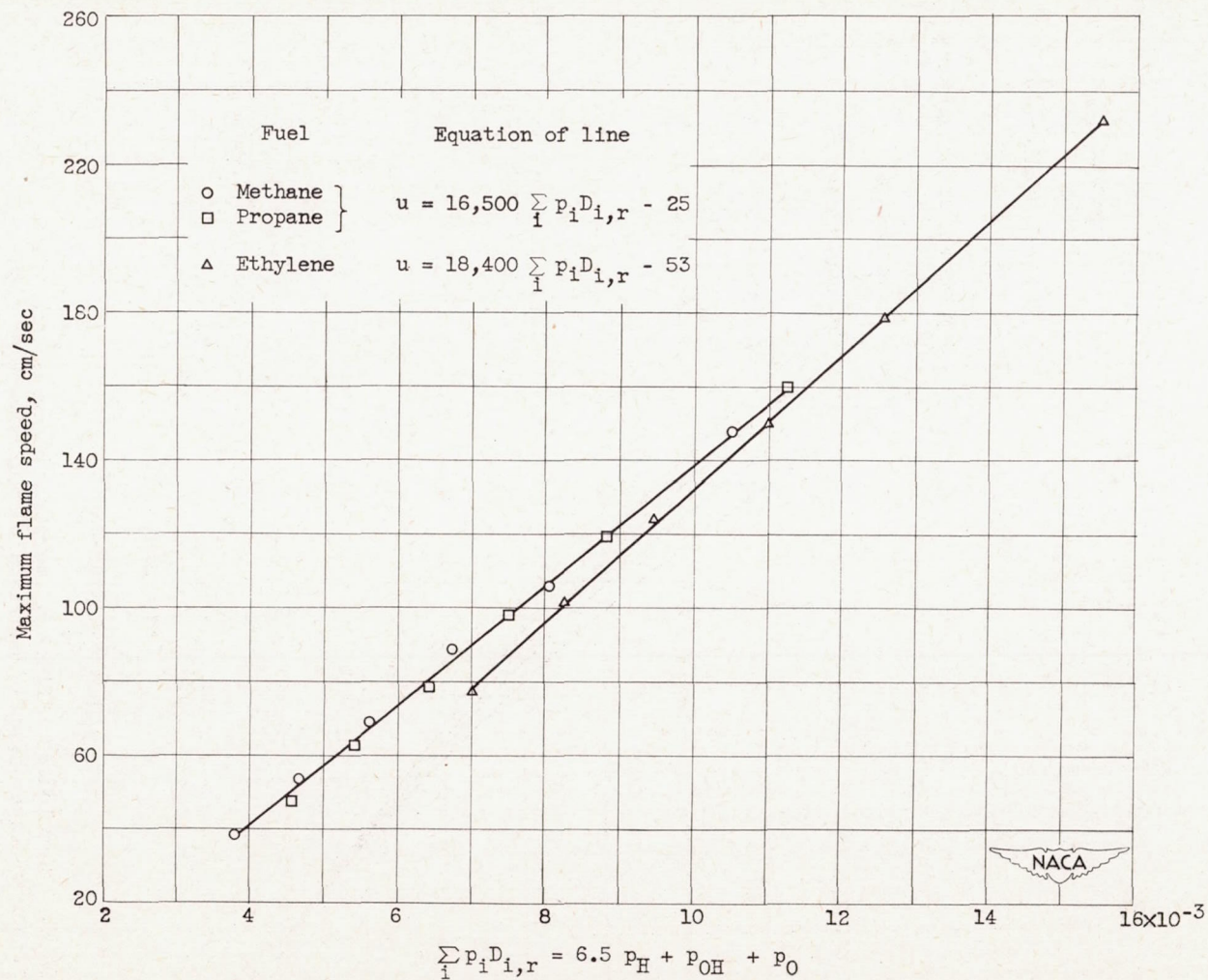
(b) Propane-air flames.

Figure 6. - Effect of temperature on maximum flame speed. Comparison of experimental results with theoretically predicted curves.



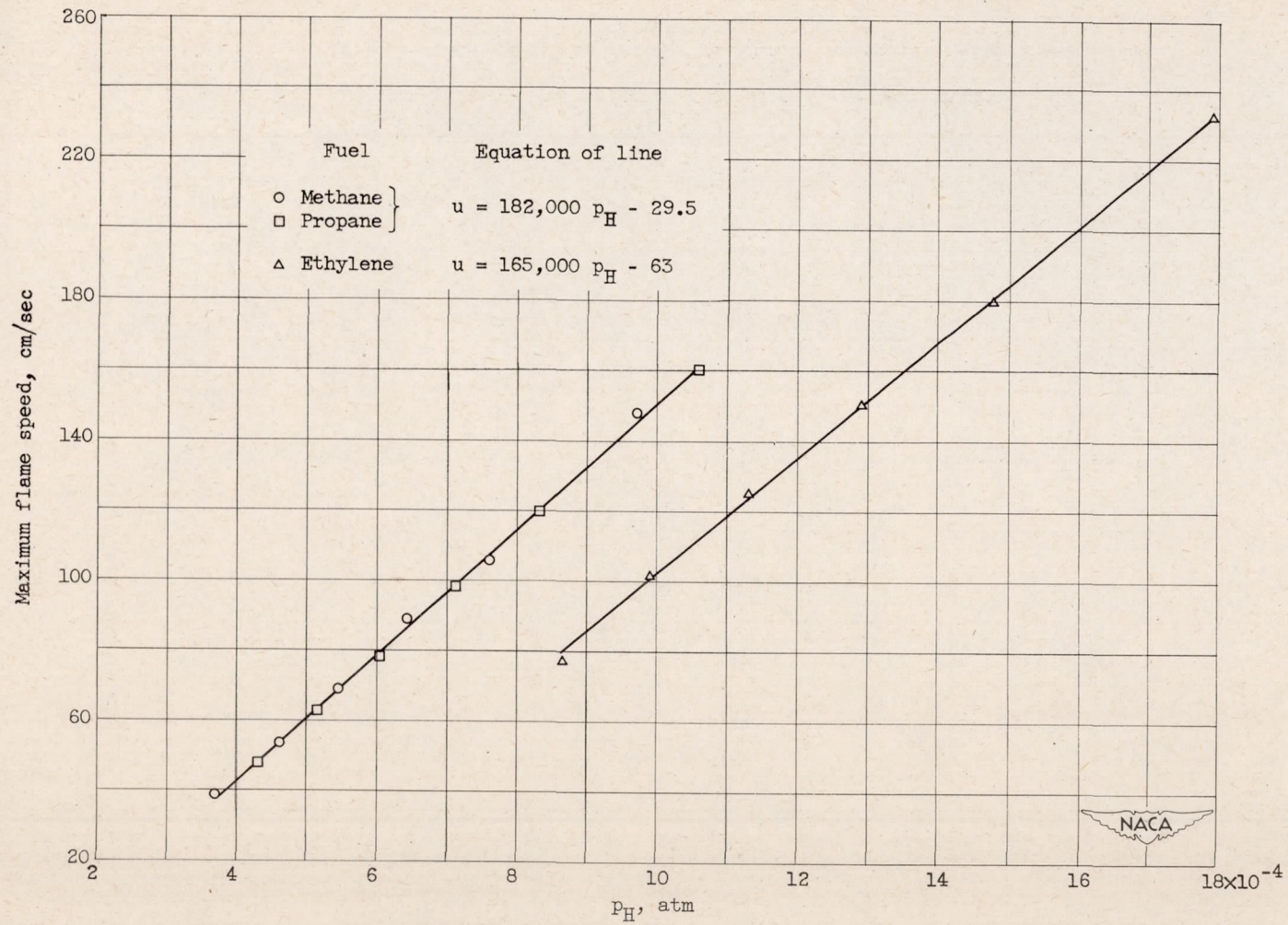
(c) Ethylene-air flames.

Figure 6. - Concluded. Effect of temperature on maximum flame speed. Comparison of experimental results with theoretically predicted curves.



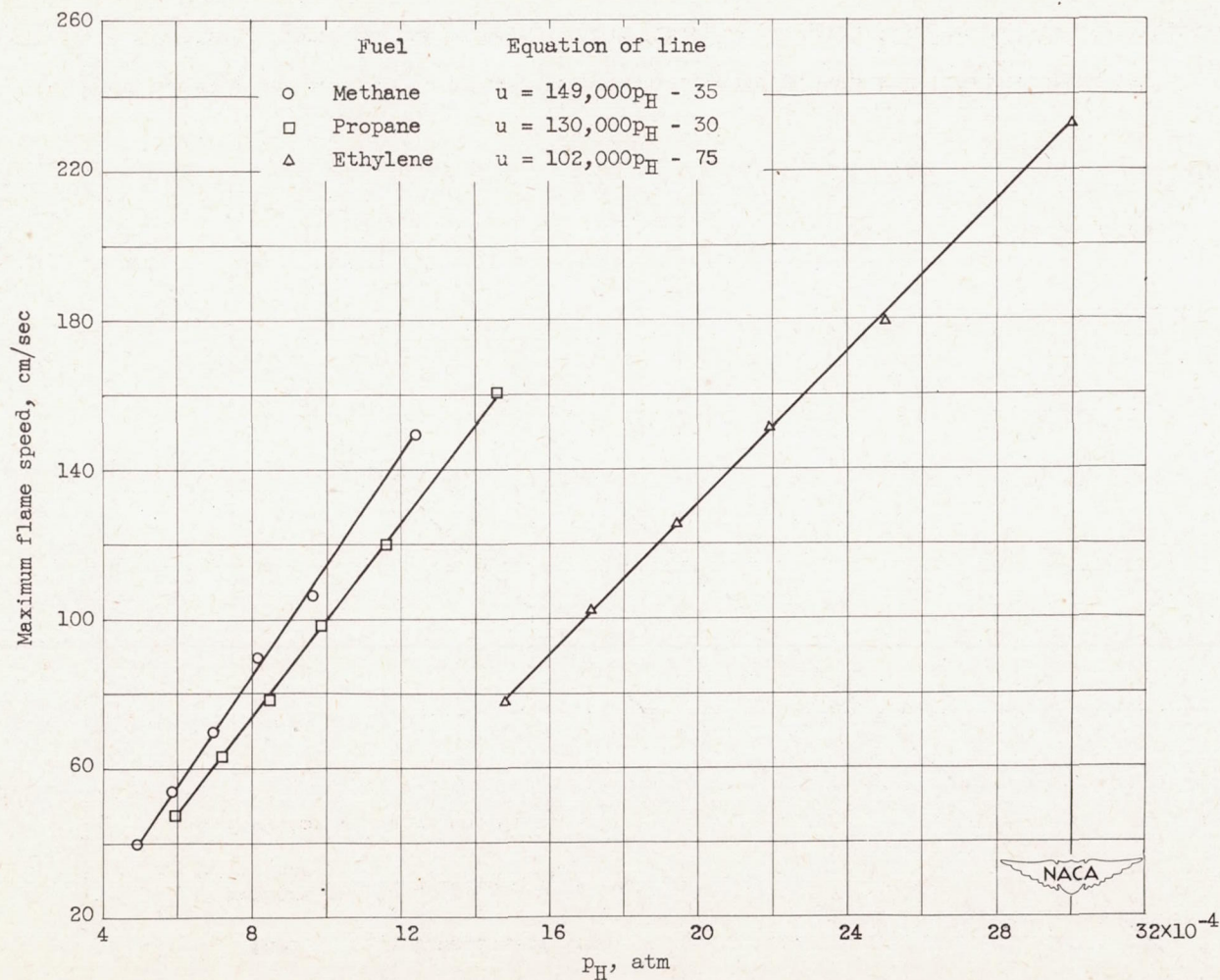
(a) Summation of effective relative radical concentrations, $6.5 p_H + p_{OH} + p_O$, at flame temperatures based on sodium D-line measurements.

Figure 7. - Correlations between maximum flame speeds and active-radical concentrations.



(b) Hydrogen-atom concentration at flame temperatures based on sodium D-line measurements.

Figure 7. - Continued. Correlations between maximum flame speeds and active-radical concentrations.



(c) Hydrogen-atom concentration at adiabatic flame temperatures.

Figure 7. - Concluded. Correlations between maximum flame speeds and active-radical concentrations.